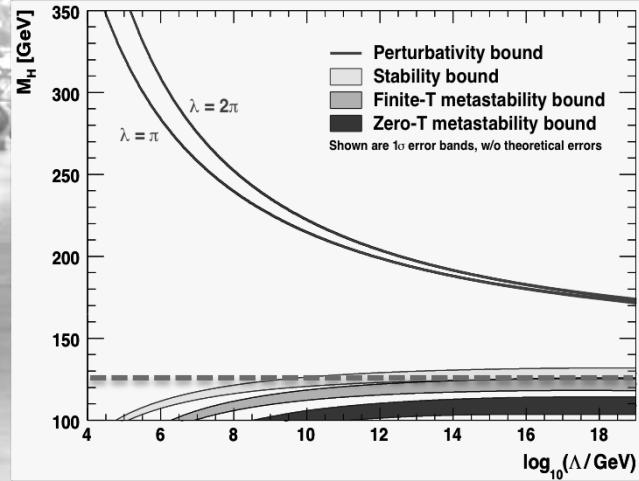
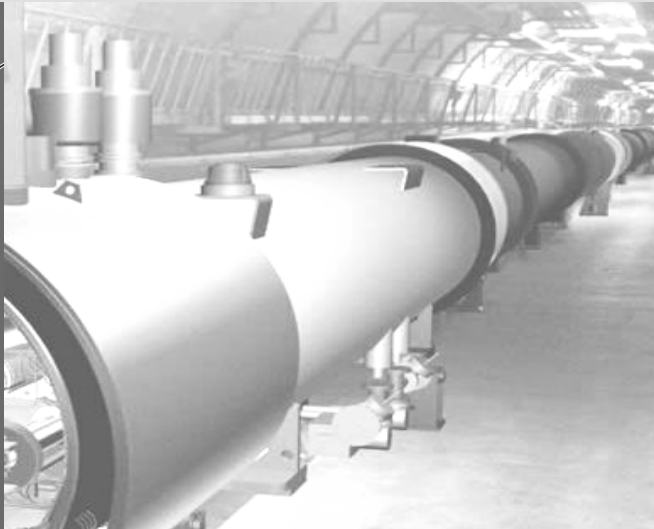
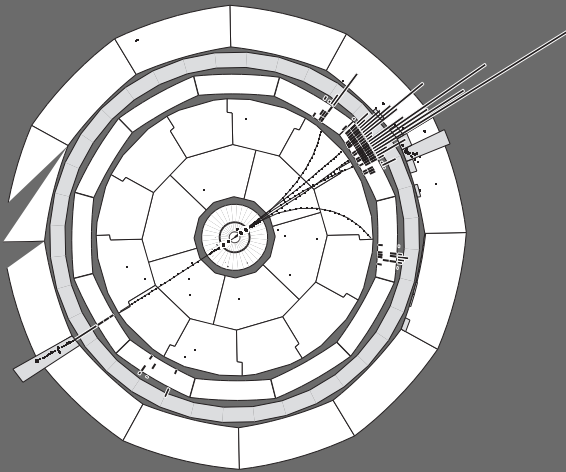
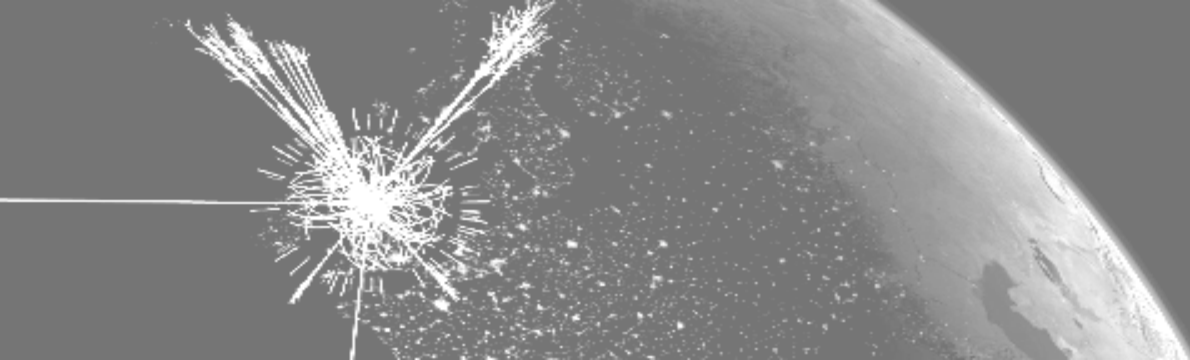


# The electroweak fit of the Standard Model after the discovery of an SM-like scalar boson

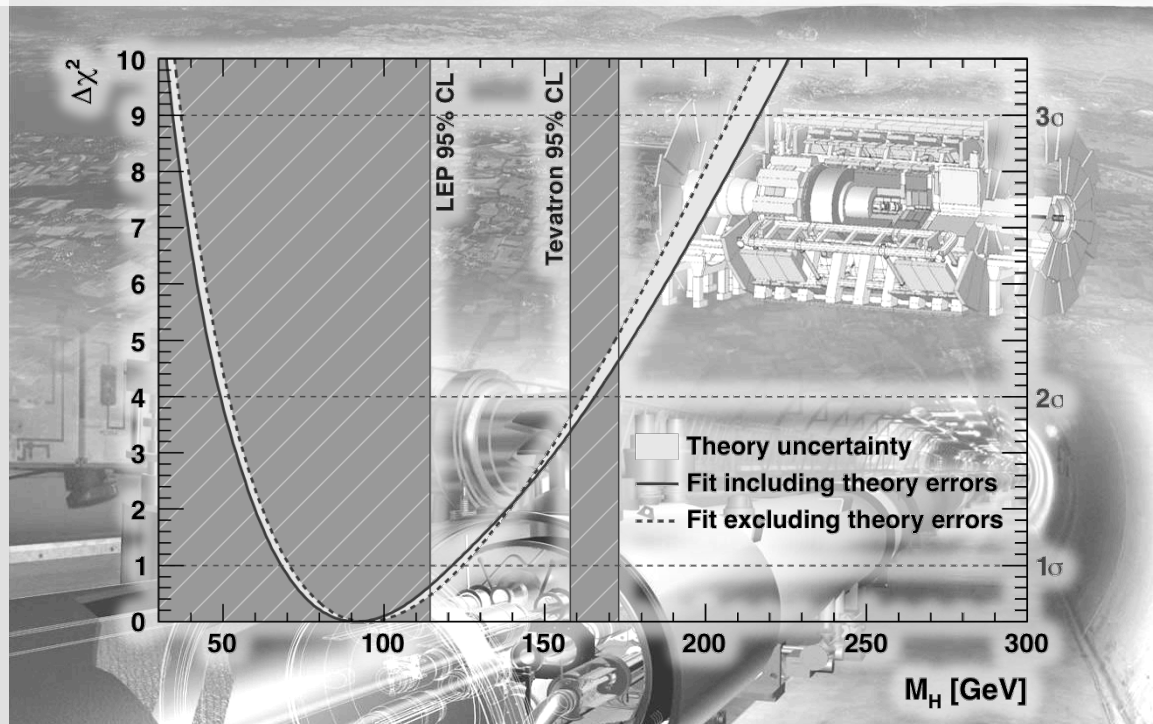
Andreas Hoecker (CERN)

Solvay workshop “Facing the Scalar Sector”, May 29-31, 2013





# Introduction



# Predictive power of the Standard Model

## Electroweak physics at the Z-pole

Vector and axial-vector couplings for  $Z \rightarrow f\bar{f}$  in SM at tree level:

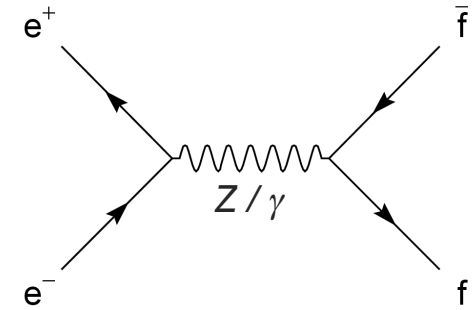
$$g_{V,f}^{(0)} \equiv g_{L,f}^{(0)} + g_{R,f}^{(0)} = I_3^f - 2Q^f \sin^2 \theta_w, \quad \sin^2 \theta_w = 1 - \frac{M_W^2}{M_Z^2}$$

$$g_{A,f}^{(0)} \equiv g_{L,f}^{(0)} - g_{R,f}^{(0)} = I_3^f$$

Electroweak unification: relation between weak and electromagnetic couplings:

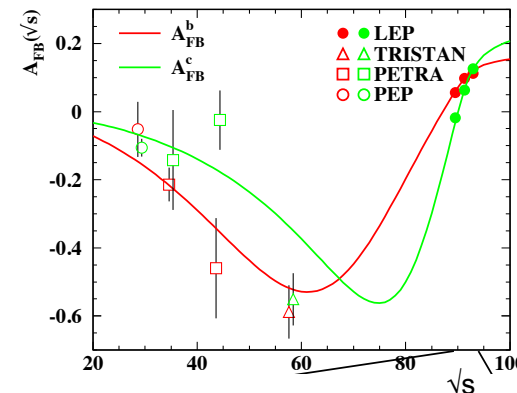
$$G_F = \frac{\pi\alpha(0)}{\sqrt{2}M_W^2(1 - M_W^2/M_Z^2)}, \quad M_W^2 = \frac{M_Z^2}{2} \cdot \left( 1 + \sqrt{1 - \frac{\sqrt{8\pi\alpha}}{G_F M_Z^2}} \right)$$

Gauge sector of SM on tree level is given by three free parameters, e.g.:  $\alpha$ ,  $M_Z$ ,  $G_F$  (best known!)



**Z-lepton coupling almost pure axial-vector**

( $\gamma$  pure vector  $\rightarrow$  large off-peak interference  $\rightarrow$  could establish Z-fermion coupling at PETRA, interesting for  $Z'$  searches via interference)



# Predictive power of the Standard Model

## Electroweak physics at the Z-pole

### Radiative corrections - modifying propagators and vertices

Significance of radiative corrections can be illustrated by verifying tree level relation:

$$\sin^2\theta_W = 1 - \frac{M_W^2}{M_Z^2}$$

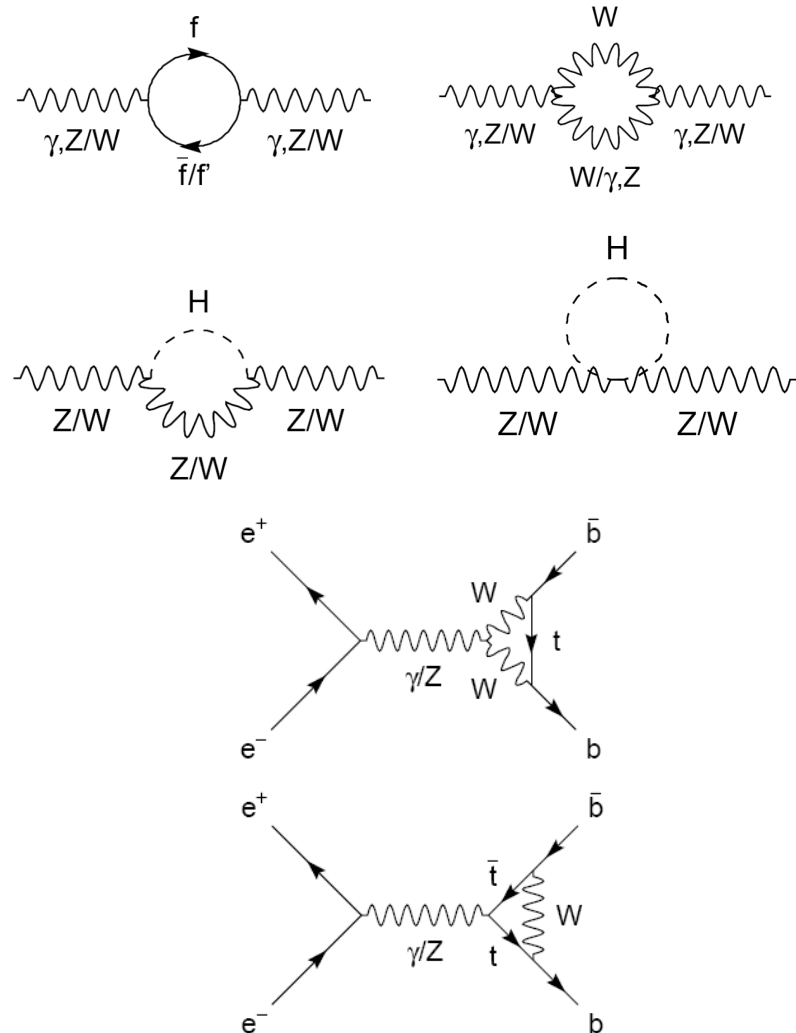
Using the measurements:

$$M_W = (80.399 \pm 0.023) \text{ GeV}$$

$$M_Z = (91.1875 \pm 0.0021) \text{ GeV}$$

one predicts:  $\sin^2\theta_W = 0.22284 \pm 0.00045$

which is  $18\sigma$  away from the experimental value obtained by combining all asymmetry measurements:  $\sin^2\theta_W = 0.23153 \pm 0.00016$



# Predictive power of the Standard Model

## Electroweak physics at the Z-pole

### Radiative corrections - modifying propagators and vertices

### Parametrisation of radiative corrections: “electroweak form-factors”: $\rho$ , $\kappa$ , $\Delta r$

- Modified (“effective”) couplings at the Z pole:

$$g_{V,f} = \sqrt{\rho_Z^f} \left( I_3^f - 2Q^f \sin^2 \theta_{\text{eff}}^f \right)$$

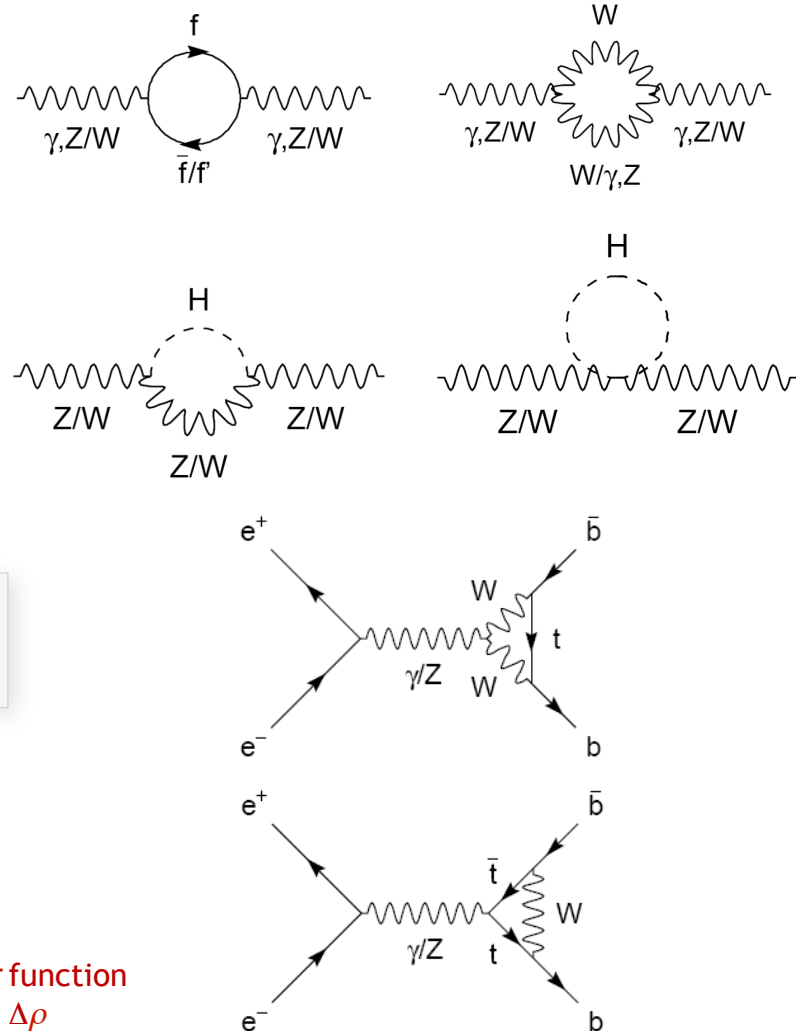
$$g_{A,f} = \sqrt{\rho_Z^f} I_3^f$$

$$\sin^2 \theta_{\text{eff}}^f = \kappa_Z^f \sin^2 \theta_W$$

$\rho$  : overall scale  
 $\kappa$  : on-shell mixing angle

- Modified W mass:

$$M_W^2 = \frac{M_Z^2}{2} \cdot \left( 1 + \sqrt{1 - \frac{\sqrt{8\pi\alpha}}{G_F M_Z^2} (1 - \Delta r)} \right) \leftarrow \Delta r \text{ function of } \Delta\rho$$



# Predictive power of the Standard Model

## Electroweak physics at the Z-pole

### Radiative corrections - modifying propagators and vertices

#### Leading order terms ( $M_W \ll M_H$ )

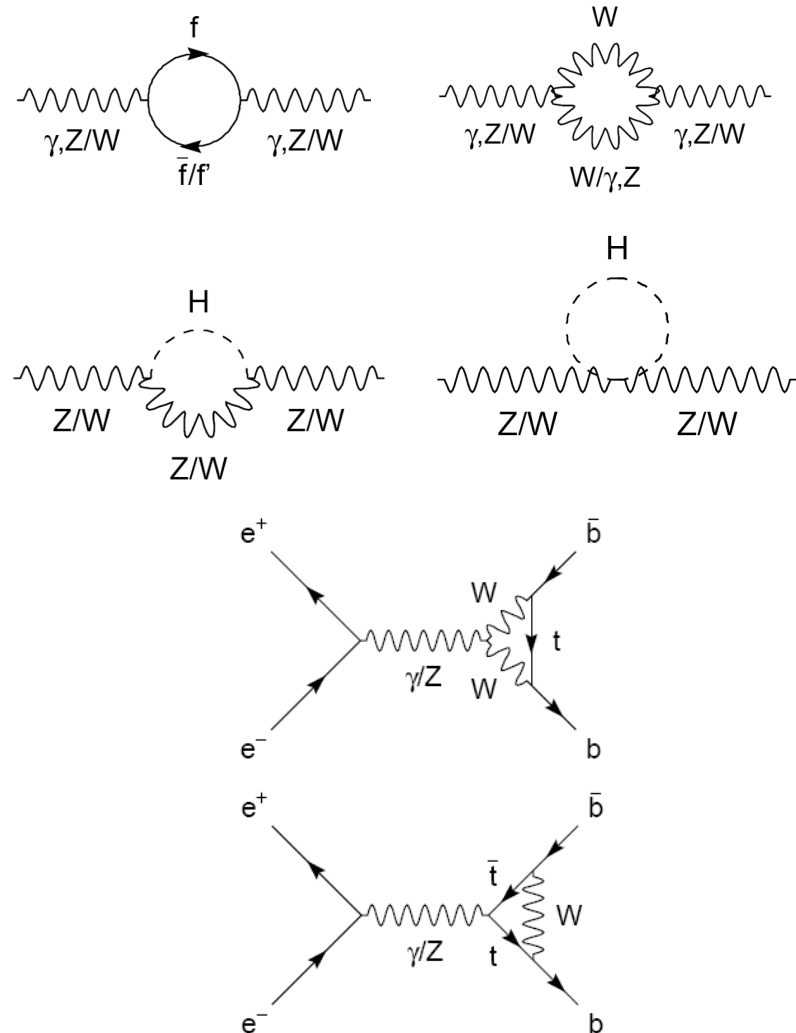
- $\rho_Z$  and  $\kappa_Z$  can be split into sum of universal contributions from propagator self-energies:

$$\Delta\rho_Z = \frac{3G_F M_W^2}{8\sqrt{2}\pi^2} \left[ \frac{m_t^2}{M_W^2} - \tan^2 \theta_w \left( \ln \frac{M_H^2}{M_W^2} - \frac{5}{6} \right) + \dots \right]$$

$$\Delta\kappa_Z = \frac{3G_F M_W^2}{8\sqrt{2}\pi^2} \left[ \frac{m_t^2}{M_W^2} \cot^2 \theta_w - \frac{10}{9} \left( \ln \frac{M_H^2}{M_W^2} - \frac{5}{6} \right) + \dots \right]$$

- and flavour-specific vertex corrections, which are very small, except for top quarks, owing to large mass and  $|V_{tb}|$  CKM element

$$\Delta\rho^f = -2\Delta\kappa^f = -\frac{G_F m_t^2}{2\sqrt{2}\pi^2} + \dots$$



# Predictive power of the Standard Model

## Electroweak physics at the Z-pole

### Radiative corrections - modifying propagators and vertices

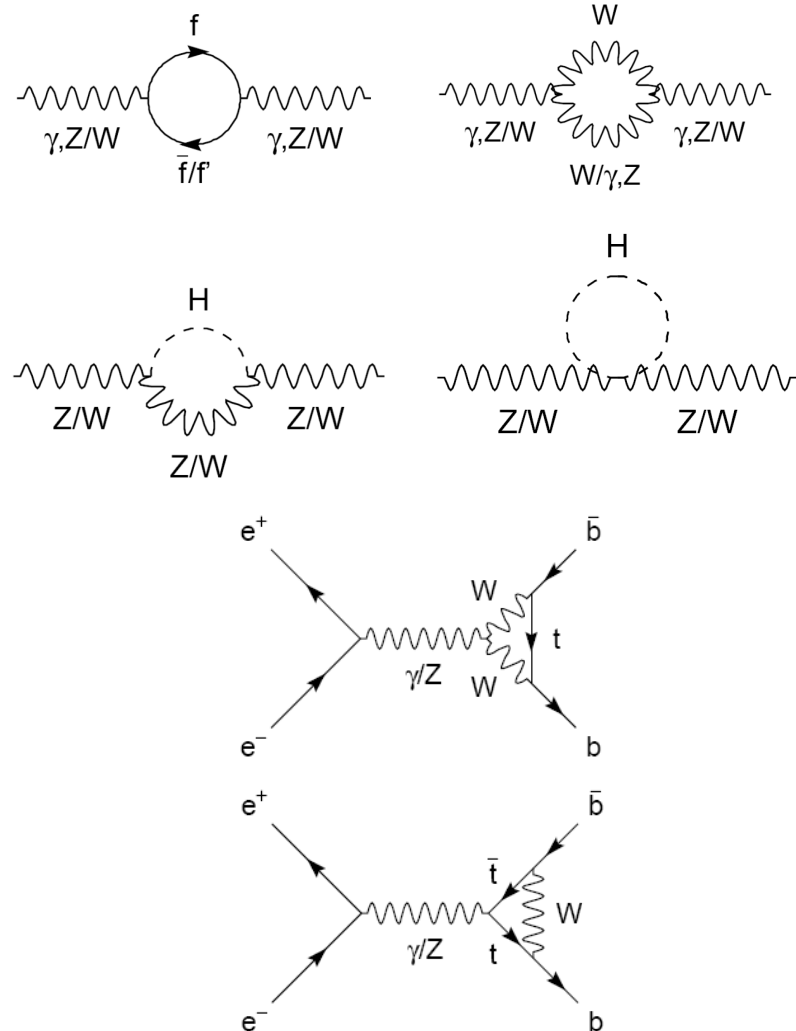
Leading order terms ( $M_W \ll M_H$ )

- $\rho_Z$  and  $\kappa_Z$  can be split into sum of universal contributions from propagator self-energies:

**Radiative corrections  
allow us to test the SM  
and to constrain unknown  
SM parameters**

- and flavour-specific vertex corrections, which are very small, except for top quarks, owing to large mass and  $|V_{tb}|$  CKM element

$$\Delta\rho^f = -2\Delta\kappa^f = -\frac{G_F m_t^2}{2\sqrt{2}\pi^2} + \dots$$



# Predictive power of the Standard Model

## Electroweak physics at the Z-pole

Observables computed using  $\rho_Z^f$ ,  $\kappa_Z^f$ ,  $\Delta r$  and QED/QCD *radiator functions*  $R_{A,f}$ ,  $R_{V,f}$

Asymmetries: 
$$A_f = \frac{2\text{Re}(g_{V,f}/g_{A,f})}{1 + \left[\text{Re}(g_{V,f}/g_{A,f})\right]^2}, \text{ where } \frac{\text{Re}(g_{V,f})}{\text{Re}(g_{A,f})} = 1 - 4|Q_f| \sin^2 \theta_{eff}^f$$

Measured asymmetries (forward-backward, left-right [+ FB] (SLD), tau polarisation) can be expressed as functions of different  $A_f$

Partial widths: 
$$\Gamma_f = N_c^f \frac{G_F M_Z^3}{6\sqrt{2}\pi} \left| \rho_Z^f \right| \left( I_3^f \right)^2 \left( \left| \frac{g_{V,f}^2}{g_{A,f}^2} \right| R_{V,f}(M_Z^2) + R_{A,f}(M_Z^2) \right)$$

Radiator functions for leptonic (hadronic) width involve QED (EW+QCD) corrections;  
 → dependence on  $\alpha_{\text{QED}}(M_Z)$  and  $\alpha_S(M_Z)$

Partial widths are highly correlated set of parameters.  
**For EW fit, use:**

- Z mass and width:  $M_Z$  ( $2 \times 10^{-5}$  accuracy!),  $\Gamma_Z$
- Hadronic pole cross section:  $\sigma_{\text{had}}^0$
- Three leptonic ratios (use lepton-univ.):  $R_\ell^0 = R_e^0 = \Gamma_{\text{had}}/\Gamma_{ee}$ ,  $R_\mu^0$ ,  $R_\tau^0$
- Hadronic width ratios:  $R_b^0 = \Gamma_{b\bar{b}}/\Gamma_{\text{had}}$ ,  $R_c^0$



# Predictive power of the Standard Model

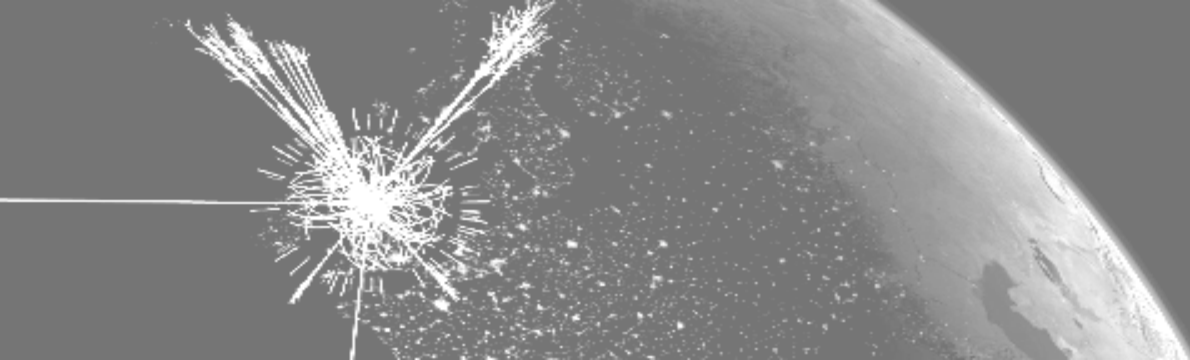
## Electroweak physics at the Z-pole

Observables computed using  $\rho_Z^f$ ,  $\kappa_Z^f$ ,  $\Delta r$  and QED/QCD *radiator functions*  $R_{A,f}$ ,  $R_{V,f}$

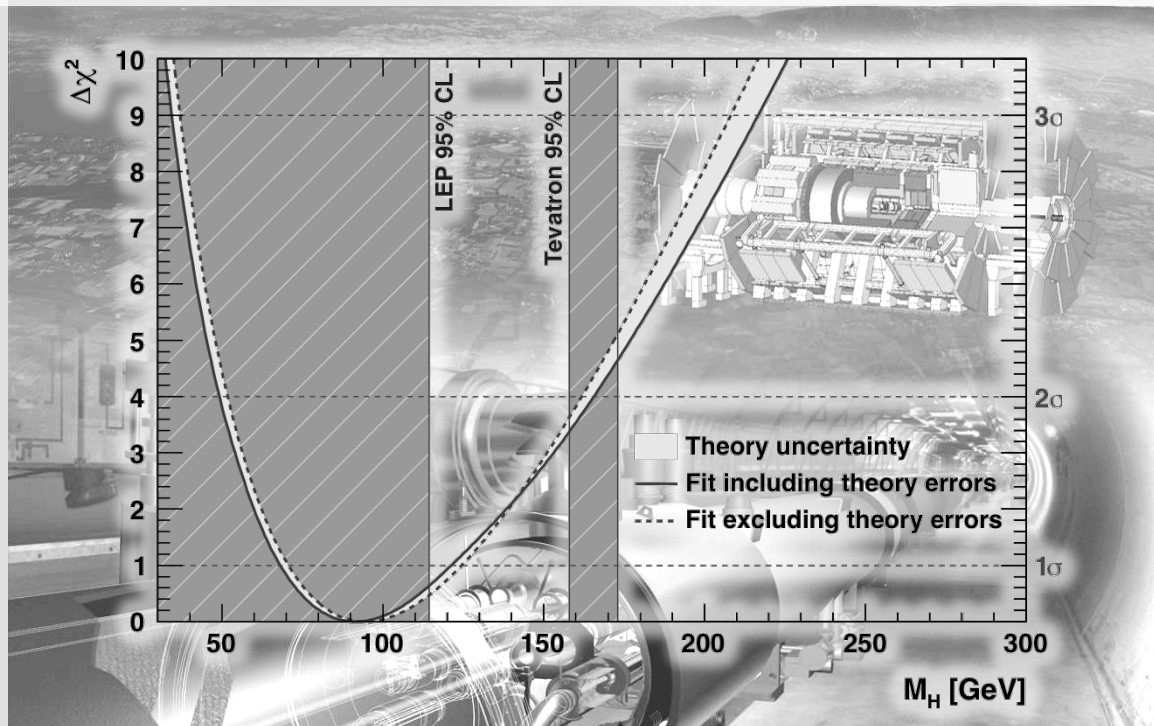
### Latest calculations for observables used

- $M_W$  **mass of the W boson**  
 $O(\alpha^2)$ ,  $O(\alpha\alpha_s)$ ,  $O(G_F\alpha_s^2 m_t^2)$ ,  $O(G_F^2\alpha_s m_t^4)$ ,  $O(G_F^3 m_t^6)$  [ Awramik et al, PRD 69, 053006 (2004)\* ]  
 $\delta_{\text{theo}} M_W = 4 \text{ MeV}$
- $\sin^2\theta_{\text{eff}}^l$  **effective weak mixing angle**  
 $O(\alpha^2)$ ,  $O(G_F^2\alpha_s m_t^4)$ ,  $O(G_F^3 m_t^6)$  [ Awramik et al, JHEP 11, 048, NP 813, 174 (2009)\* ]  
 $\delta_{\text{theo}} \sin^2\theta_{\text{eff}}^l = 4.7 \times 10^{-5}$
- $\Gamma_Z$ ,  $\Gamma_W$  **Total widths of Z and W** [ Cho et al, arXiv:1104\* ]
- $R_l$  **leptonic width ratio**  
QCD Adler functions at 3NLO [ Baikov et al., PRL 108, 222003 (2012)\* ]  
 $\alpha_{\text{QED}}(M_Z)$  from newest hadronic data [ Davier et al., EPJ.C71, 1515 (2011) ]
- $R_b$   **$Z \rightarrow bb$  width ratio**  
Full two-loop fermionic correction (sizable: theoretical uncertainties larger than expected?) [ Freitas et al, JHEP 08, 050 (2012)\* ]

*\*References only those used directly by Gfitter.  
Full list of theoretical calculations referenced given in 0811.0009.*



# The electroweak fit



# Electroweak fits

Several groups perform these fits with regular updates (LEPEWWG, PDG, Gfitter, BSM groups)

## A long tradition

- Precision measurements crucial. After LEP/SLC era, results from Tevatron & soon also LHC
- Huge & pioneering work to compute loop corrections to two-loop order or higher

## Brout-Englert-Higgs hunting

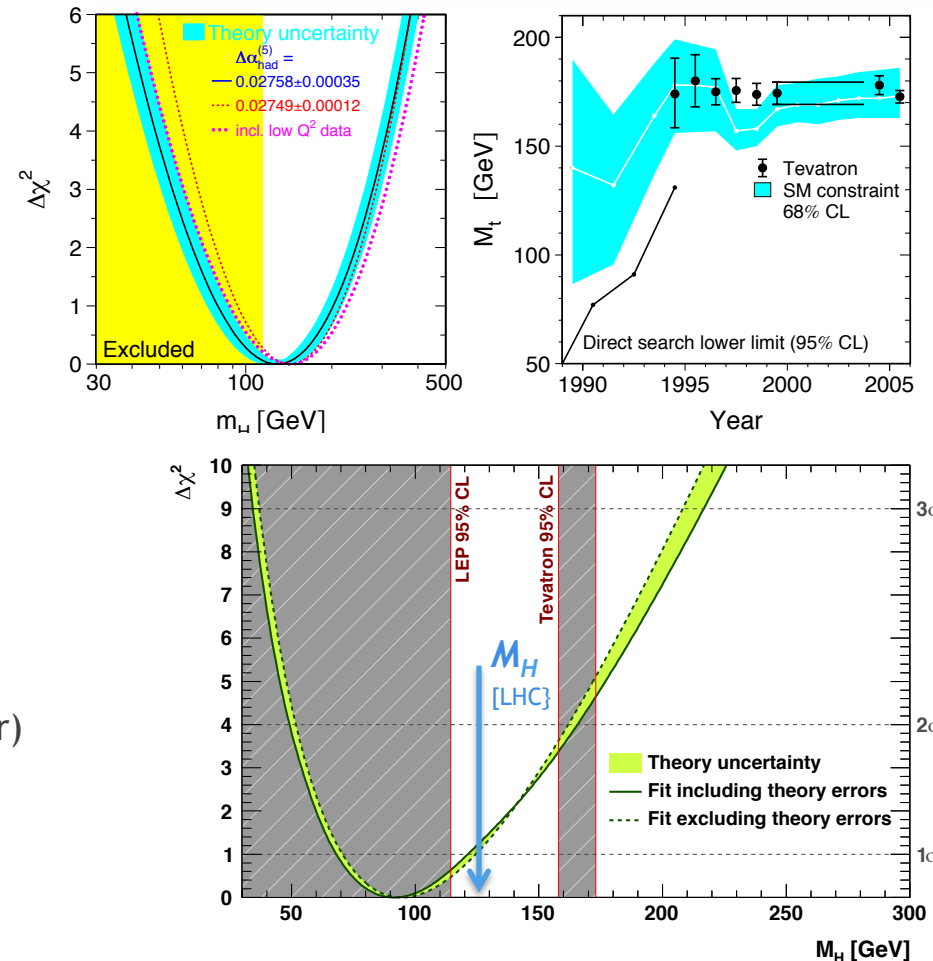
- $M_H$  last missing parameter of the SM
- Indirect determination (2011):  $M_H = 96^{+31}_{-24}$  GeV
- Exclusion limits were incorporated in EW fits

## Discovery of new boson in July 2012

- The cross section and branching ratios are (so far) compatible with the SM scalar boson
- Assume in the following that the boson is the SM scalar:  $M_H = 125.7 \pm 0.4$  GeV\*

\*Exact value and uncertainty irrelevant for EW fit in SM

LEP: Phys. Rept. 427 (2006) 257, Gfitter: EPJ C72 (2012) 2003



# Experimental observables

Several groups perform these fits with regular updates (LEPEWWG, PDG, Gfitter, BSM groups)

## Experimental inputs:

- **Z-pole observables:** LEP/SLD results (corrected for ISR/FSR QED effects)  
[ADLO & SLD, Phys. Rept. 427, 257 (2006)]
  - **Total and partial cross sections** around Z:  $M_Z, \Gamma_Z, \sigma_{\text{had}}^0, R_l^0, R_c^0, R_b^0$
  - **Asymmetries** on the Z pole:  $A_{\text{FB}}^{0,l}, A_{\text{FB}}^{0,b}, A_{\text{FB}}^{0,c}, A_l, A_c, A_b, \sin^2\theta_{\text{eff}}^l (Q_{\text{FB}})$
- $M_W$  and  $\Gamma_W$ : LEP + Tevatron average [arXiv:1204:0042]
- $m_t$ : latest Tevatron average [CDF & D0, new combination, arXiv:1305.3929]
- $m_c, m_b$ : world averages [PDG, Phys. Lett. B667, 1 (2008) and 2009 partial update for the 2010 edition]
- $\Delta\alpha_{\text{had}}(M_Z)$ : data + QCD-driven [Davier et al., EPJ.C71, 1515 (2011) + rescaling mechanism to account for  $\alpha_s$  dependency]
- $M_H$ : LHC [arXiv:1207.7214, arXiv:1207.7235]

## Fit parameters

- $\Delta\alpha_{\text{had}}(M_Z), \alpha_S(M_Z), M_Z, M_H, m_c, m_b, m_t$  + theory uncertainty parameters  $\delta_{\text{theo}} M_W / \sin^2\theta_{\text{eff}}^l$
- Other parameters well enough known and fixed in fit

Parameter	Input value	
$M_Z$ [GeV]	$91.1875 \pm 0.0021$	LEP
$\Gamma_Z$ [GeV]	$2.4952 \pm 0.0023$	
$\sigma_{\text{had}}^0$ [nb]	$41.540 \pm 0.037$	
$R_\ell^0$	$20.767 \pm 0.025$	
$A_{\text{FB}}^{0,\ell}$	$0.0171 \pm 0.0010$	SLC
$A_\ell^{(*)}$	$0.1499 \pm 0.0018$	
$\sin^2\theta_{\text{eff}}^\ell(Q_{\text{FB}})$	$0.2324 \pm 0.0012$	SLC
$A_c$	$0.670 \pm 0.027$	
$A_b$	$0.923 \pm 0.020$	LEP
$A_{\text{FB}}^{0,c}$	$0.0707 \pm 0.0035$	
$A_{\text{FB}}^{0,b}$	$0.0992 \pm 0.0016$	SLC
$R_c^0$	$0.1721 \pm 0.0030$	
$R_b^0$	$0.21629 \pm 0.00066$	

Parameter	Input value	
$M_H$ [GeV] <sup>(<math>\circ</math>)</sup>	$125.7 \pm 0.4$	LHC
$M_W$ [GeV]	$80.385 \pm 0.015$	Tevatron & LEP
$\Gamma_W$ [GeV]	$2.085 \pm 0.042$	
$\bar{m}_c$ [GeV]	$1.27^{+0.07}_{-0.11}$	Tevatron
$\bar{m}_b$ [GeV]	$4.20^{+0.17}_{-0.07}$	
$m_t$ [GeV]	$173.20 \pm 0.87$	Tevatron
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$ ( $\Delta\nabla$ )	$2757 \pm 10$	
$\alpha_s(M_Z^2)$	–	
$\delta_{\text{th}} M_W$ [MeV]	$[-4, 4]_{\text{theo}}$	
$\delta_{\text{th}} \sin^2\theta_{\text{eff}}^\ell$ ( $\Delta$ )	$[-4.7, 4.7]_{\text{theo}}$	

Correlations for observables from Z lineshape fit

	$M_Z$	$\Gamma_Z$	$\sigma_{\text{had}}^0$	$R_\ell^0$	$A_{\text{FB}}^{0,\ell}$
$M_Z$	1	-0.02	-0.05	0.03	0.06
$\Gamma_Z$		1	-0.30	0.00	0.00
$\sigma_{\text{had}}^0$			1	0.18	0.01
$R_\ell^0$				1	-0.06
$A_{\text{FB}}^{0,\ell}$					1

Correlations for heavy-flavour observables at Z pole

	$A_{\text{FB}}^{0,c}$	$A_{\text{FB}}^{0,b}$	$A_c$	$A_b$	$R_c^0$	$R_b^0$
$A_{\text{FB}}^{0,c}$	1	0.15	0.04	-0.02	-0.06	0.07
$A_{\text{FB}}^{0,b}$		1	0.01	0.06	0.04	-0.10
$A_c$			1	0.11	-0.06	0.04
$A_b$				1	0.04	-0.08
$R_c^0$					1	-0.18

Parameter	Input value	Free in fit	Fit Result	Fit without $M_H$ measurements	Fit without exp. input in line
$M_H$ [GeV] <sup>o</sup>	$125.7 \pm 0.4$	yes	$125.7 \pm 0.4$	$94.1^{+25}_{-22}$	$94.1^{+25}_{-22}$
$M_W$ [GeV]	$80.385 \pm 0.015$	–	$80.367^{+0.006}_{-0.007}$	$80.380^{+0.011}_{-0.012}$	$80.360 \pm 0.011$
$\Gamma_W$ [GeV]	$2.085 \pm 0.042$	–	$2.091 \pm 0.001$	$2.092 \pm 0.001$	$2.091 \pm 0.001$
$M_Z$ [GeV]	$91.1875 \pm 0.0021$	yes	$91.1878 \pm 0.0021$	$91.1874 \pm 0.0021$	$91.1983 \pm 0.0115$
$\Gamma_Z$ [GeV]	$2.4952 \pm 0.0023$	–	$2.4953 \pm 0.0014$	$2.4957 \pm 0.0015$	$2.4949 \pm 0.0017$
$\sigma_{\text{had}}^0$ [nb]	$41.540 \pm 0.037$	–	$41.480 \pm 0.014$	$41.479 \pm 0.014$	$41.472 \pm 0.015$
$R_\ell^0$	$20.767 \pm 0.025$	–	$20.739 \pm 0.017$	$20.741 \pm 0.017$	$20.713 \pm 0.026$
$A_{\text{FB}}^{0,\ell}$	$0.0171 \pm 0.0010$	–	$0.01627^{+0.0001}_{-0.0002}$	$0.01637 \pm 0.0002$	$0.01624 \pm 0.0002$
$A_\ell$ (*)	$0.1499 \pm 0.0018$	–	$0.1473^{+0.0006}_{-0.0008}$	$0.1477^{+0.0009}_{-0.0008}$	–
$\sin^2\theta_{\text{eff}}^\ell(Q_{\text{FB}})$	$0.2324 \pm 0.0012$	–	$0.23148^{+0.00011}_{-0.00007}$	$0.23143^{+0.00010}_{-0.00012}$	$0.23150 \pm 0.00009$
$A_c$	$0.670 \pm 0.027$	–	$0.6681^{+0.00021}_{-0.00042}$	$0.6682^{+0.00042}_{-0.00035}$	$0.6680 \pm 0.00031$
$A_b$	$0.923 \pm 0.020$	–	$0.93464^{+0.00005}_{-0.00007}$	$0.93468^{+0.00008}_{-0.00007}$	$0.93463 \pm 0.00006$
$A_{\text{FB}}^{0,c}$	$0.0707 \pm 0.0035$	–	$0.0739^{+0.0003}_{-0.0005}$	$0.0740^{+0.0005}_{-0.0004}$	$0.0738 \pm 0.0004$
$A_{\text{FB}}^{0,b}$	$0.0992 \pm 0.0016$	–	$0.1032^{+0.0004}_{-0.0006}$	$0.1036^{+0.0007}_{-0.0006}$	$0.1034 \pm 0.0003$
$R_c^0$	$0.1721 \pm 0.0030$	–	$0.17222^{+0.00006}_{-0.00005}$	$0.17223 \pm 0.00006$	$0.17223 \pm 0.00006$
$R_b^0$	$0.21629 \pm 0.00066$	–	$0.21491 \pm 0.00005$	$0.21492 \pm 0.00005$	$0.21490 \pm 0.00005$
$\overline{m}_c$ [GeV]	$1.27^{+0.07}_{-0.11}$	yes	$1.27^{+0.07}_{-0.11}$	$1.27^{+0.07}_{-0.11}$	–
$\overline{m}_b$ [GeV]	$4.20^{+0.17}_{-0.07}$	yes	$4.20^{+0.17}_{-0.07}$	$4.20^{+0.17}_{-0.07}$	–
$m_t$ [GeV]	$173.20 \pm 0.87$	yes	$173.49 \pm 0.82$	$173.17 \pm 0.86$	$175.83^{+2.74}_{-2.42}$
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$ († $\Delta$ )	$2756 \pm 10$	yes	$2755 \pm 11$	$2757 \pm 11$	$2716^{+49}_{-43}$
$\alpha_s(M_Z^2)$	–	yes	$0.1188^{+0.0028}_{-0.0027}$	$0.1190^{+0.0028}_{-0.0027}$	$0.1188 \pm 0.0027$
$\delta_{\text{th}} M_W$ [MeV]	$[-4, 4]_{\text{theo}}$	yes	4	4	–
$\delta_{\text{th}} \sin^2\theta_{\text{eff}}^\ell$ (†)	$[-4.7, 4.7]_{\text{theo}}$	yes	–1.4	4.7	–

(<sup>o</sup>) Average of ATLAS ( $M_H = 126.0 \pm 0.4$  (stat)  $\pm 0.4$  (sys)) and CMS ( $M_H = 125.3 \pm 0.4$  (stat)  $\pm 0.5$  (sys)) measurements assuming no correlation of the systematic uncertainties. (\*) Average of LEP ( $A_\ell = 0.1465 \pm 0.0033$ ) and SLD ( $A_\ell = 0.1513 \pm 0.0021$ ) measurements, used as two measurements in the fit. The fit w/o the LEP (SLD) measurement gives  $A_\ell = 0.1474^{+0.0005}_{-0.0009}$  ( $A_\ell = 0.1467^{+0.0006}_{-0.0004}$ ). (†) In units of  $10^{-5}$ . ( $\Delta$ ) Rescaled due to  $\alpha_s$  dependency.

## Goodness-of-fit:

$$\chi^2_{\min}/n_{\text{dof}} = 20.7/14 \rightarrow p\text{-value} = 0.09_{\text{toy-MC}}$$

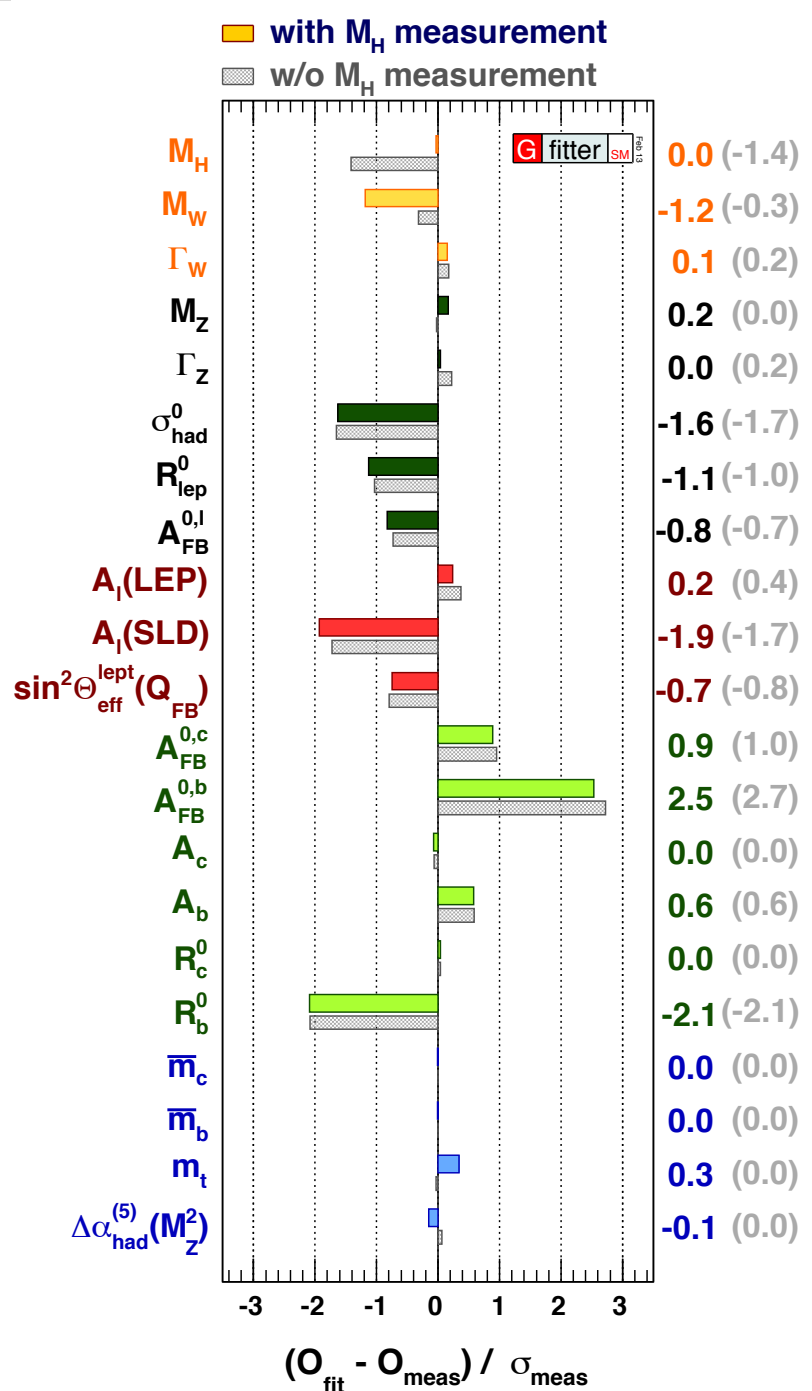
- ▶ large value of  $\chi^2_{\min}$  not due to  $M_H$  measurement

- ▶ Without  $M_H$  measurement:

$$\chi^2_{\min}/n_{\text{dof}} = 19.3/13 \rightarrow p\text{-value} \sim 0.11$$

## Pull values after fit:

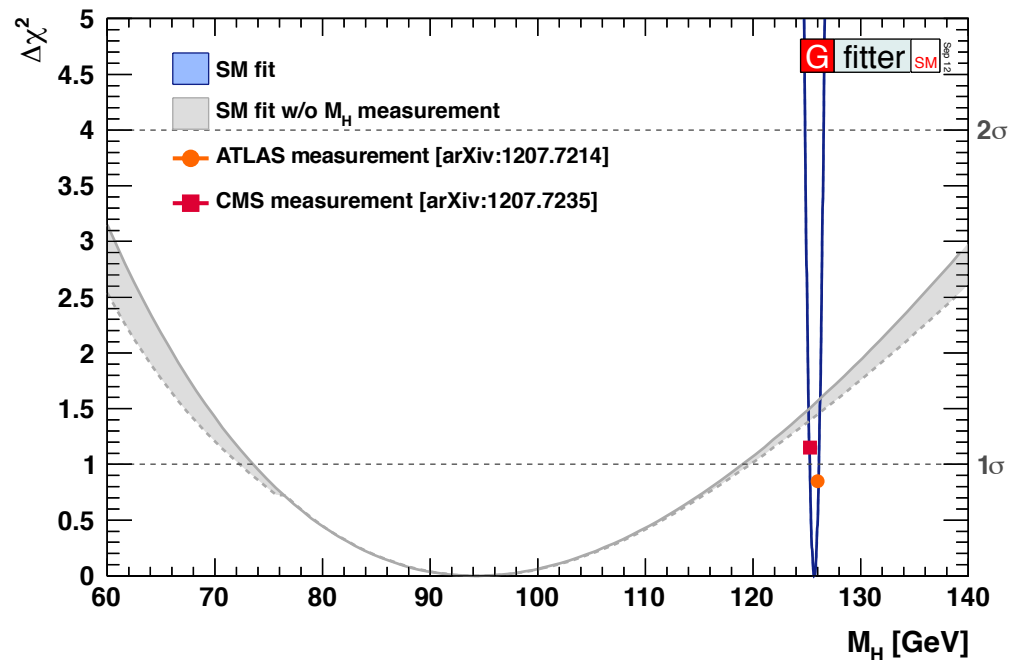
- ▶ No pull value exceeds deviation of more than  $3\sigma$  (consistency of SM)
- ▶ Small pulls for  $M_H$ ,  $A_c$ ,  $R_c^0$ ,  $m_c$  and  $m_b$  indicate that their input accuracies exceed the fit requirements
- ▶ Largest pulls in  $b$ -sector:  $A^{0,b}_{\text{FB}}$  and  $R^0_b$  with  $2.5\sigma$  and  $-2.1\sigma$  (little dependence on  $M_H$ )
- ▶ For comparison:  $R^0_b$  using one fermionic loop calculation:  $0.8\sigma$



Plot inspired by Eberhardt et al. [arXiv:1209.1101]

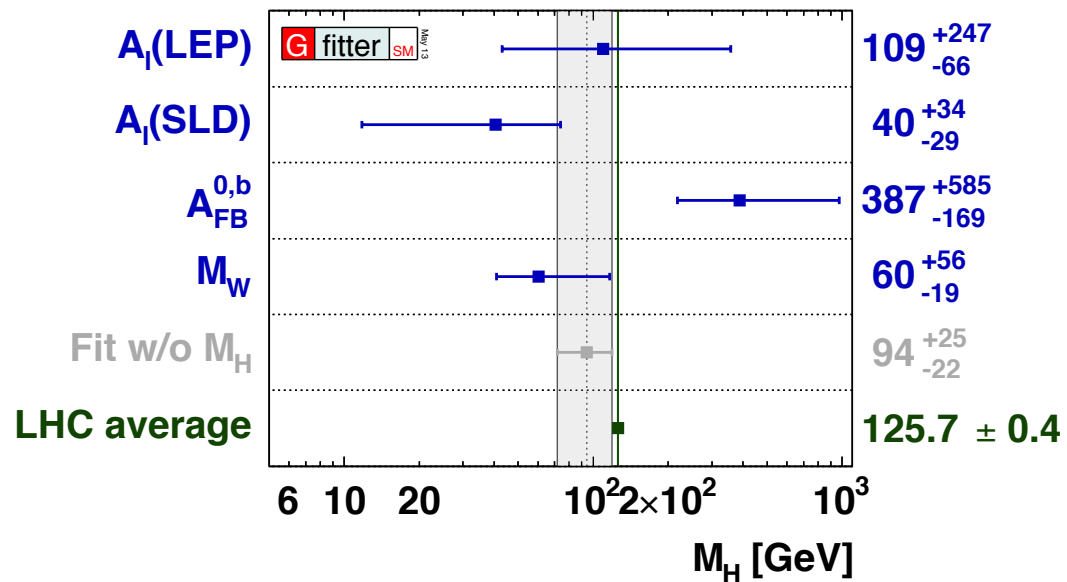
## Scan of the $\Delta\chi^2$ profile versus $M_H$

- ▶ Blue line: full SM fit
- ▶ Grey band: fit without  $M_H$  measurement
- ▶ Fit without MH input gives  $M_H = 94^{+25}_{-22}$  GeV
- ▶ Consistent within  $1.3\sigma$  with measurement



## Tension in $M_H$ fit ?

- ▶ Determination of  $M_H$  removing all sensitive observables except the given one
- ▶ Tension ( $2.5\sigma$  - from toy MC) between  $A_{FB}^{0,b}$ ,  $A_l(\text{SLD})$  and  $M_W$





# Global fit results

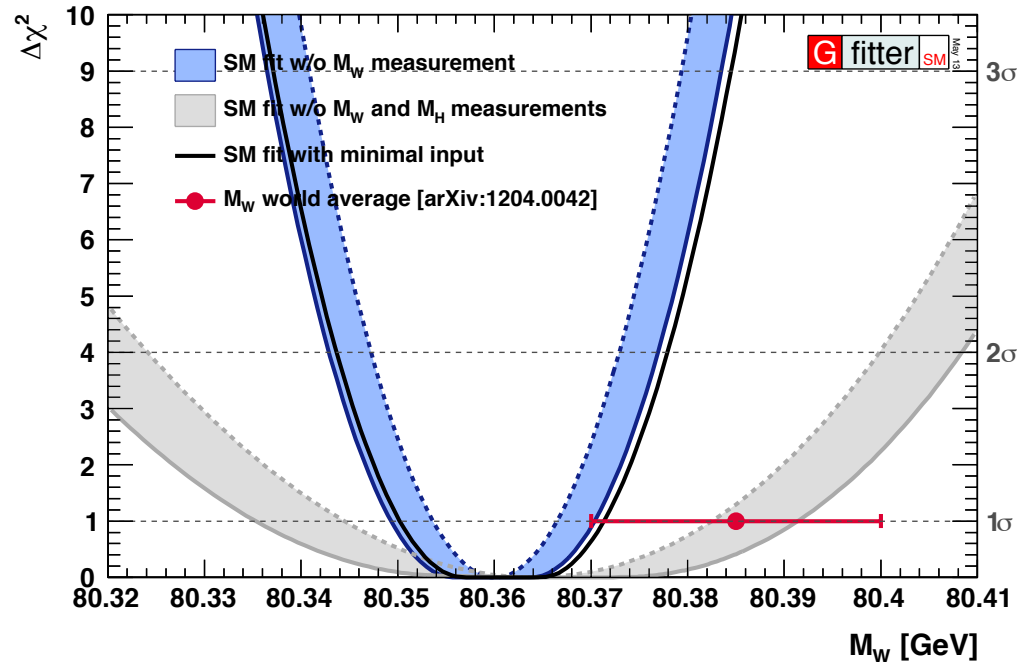
## Indirect determination of the $W$ boson mass

### Scan of $\Delta\chi^2$ profile versus $M_W$

- ▶  $M_H$  measurement allows for precise constraint of  $M_W$
- ▶ Also shown SM fit with minimal input:  $M_Z$ ,  $G_F$ ,  $\Delta\alpha_{\text{had}}(M_Z)$ ,  $\alpha(M_Z)$ ,  $M_H$  and fermion masses
- ▶ Consistency between total fit and SM fit with minimal input

Fit results in the indirect determination :

$$\begin{aligned} M_W &= 80.3603 \pm 0.0056(m_{\text{top}}) \pm 0.0026(M_Z) \pm 0.0018(\Delta\alpha_{\text{had}}) \\ &\quad \pm 0.0027(\alpha_s) \pm 0.0002(M_H) \pm 0.0040(\text{theo}) \text{ GeV} \\ &= 80.360 \pm 0.011(\text{tot}) \text{ GeV, more precise than experimental value} \\ &= 80.385 \pm 0.015(\text{exp}) \text{ GeV} \quad [\text{Tevatron/LEP: arXiv:1204.0042}] \end{aligned}$$



# Global fit results

## Effective weak mixing angle

### Scan of $\Delta\chi^2$ profile versus $\sin^2\theta_{\text{eff}}^l$

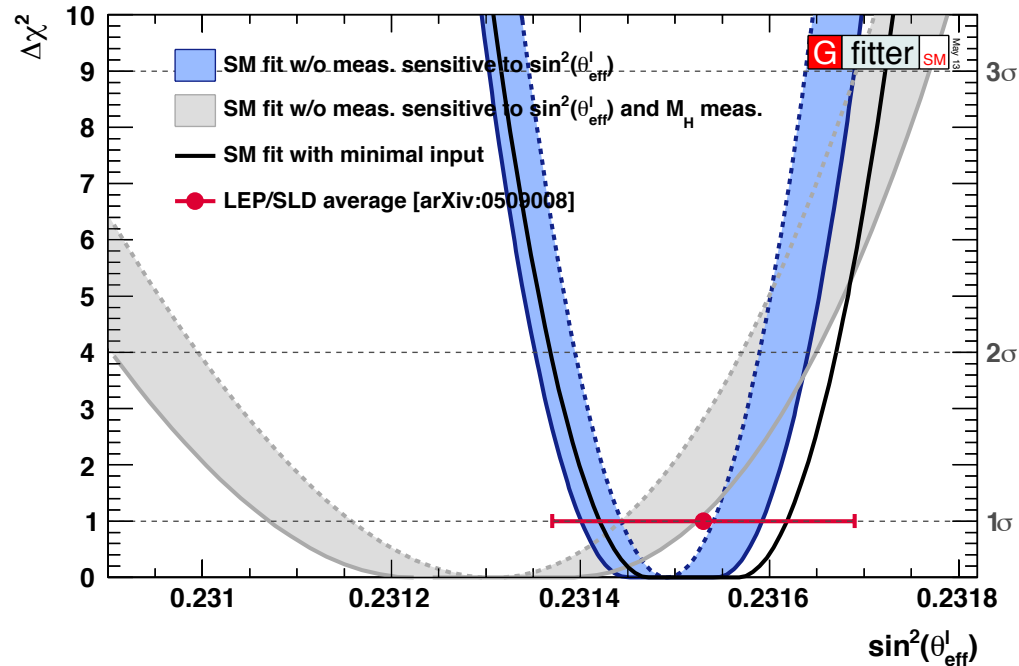
- ▶ All observables sensitive to  $\sin^2\theta_{\text{eff}}^l$  removed from fit
- ▶  $M_H$  measurement allows for precise constraint
- ▶ Also shown SM fit with minimal input

Fit results in the indirect determination :

$$\sin^2\theta_{\text{eff}}^l = 0.231496 \pm 0.000030(m_{\text{top}}) \pm 0.000015(M_Z) \pm 0.000035(\Delta\alpha_{\text{had}}) \\ \pm 0.000010(\alpha_S) \pm 0.000002(M_H) \pm 0.000047(\text{theo})$$

$$= 0.23150 \pm 0.00010(\text{tot}), \text{ more precise than LEP/SLD average}$$

$$= 0.23153 \pm 0.00016(\text{exp}) \quad [\text{LEP/SLD: Phys Rept 427 (2006) 257}]$$



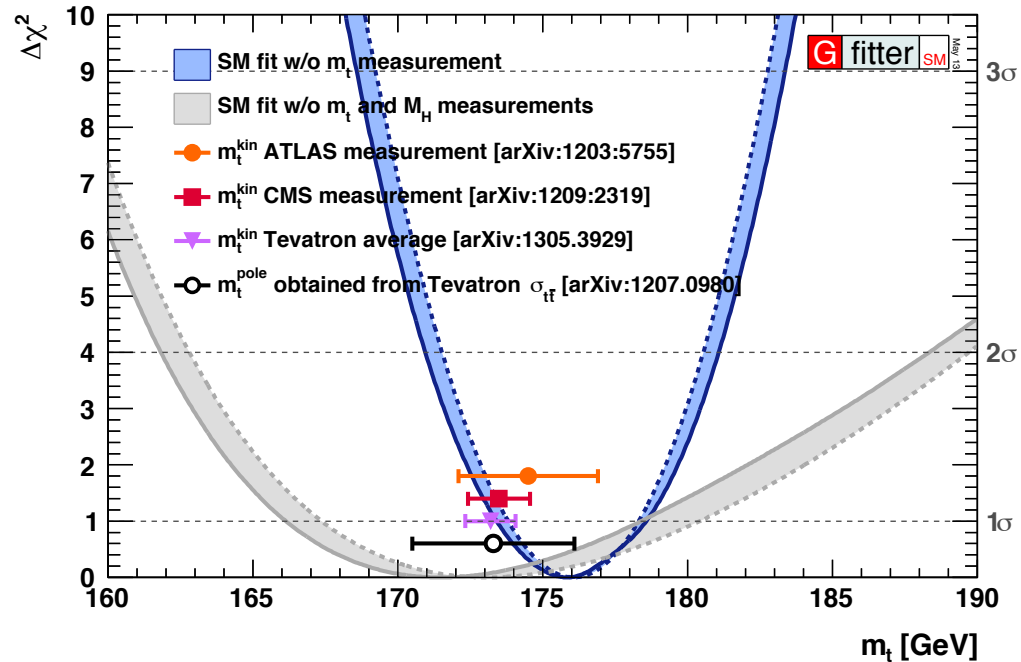
# Global fit results

## Top mass

### Scan of $\Delta\chi^2$ profile versus $m_{\text{top}}$

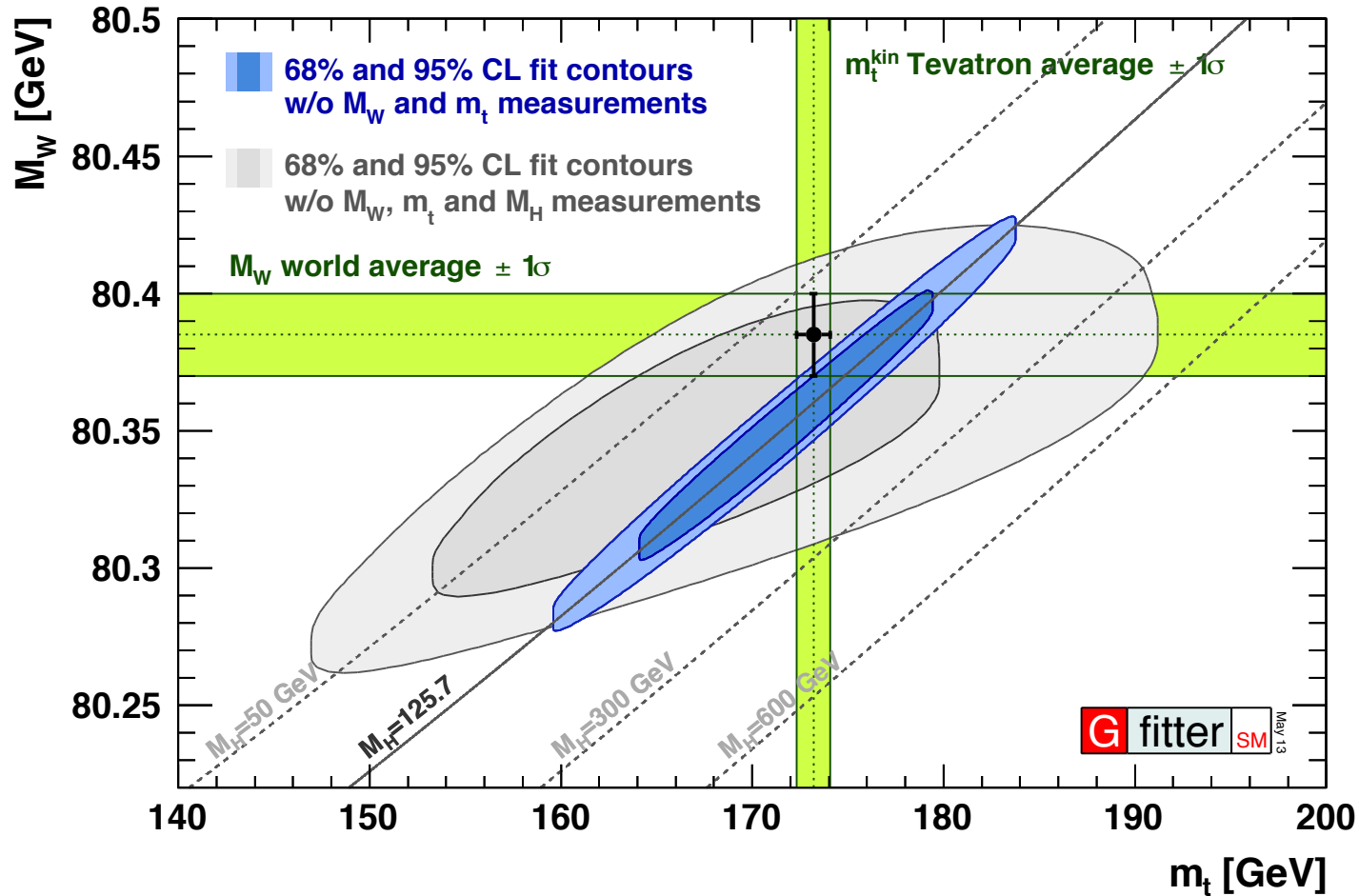
Fit results in the indirect determination :

$$\begin{aligned} m_{\text{top}} &= 175.8^{+2.7}_{-2.4} \text{ (tot) GeV} \\ &= 173.2 \pm 0.9 \text{ (exp) [Tevatron: arXiv:1207.0980]} \end{aligned}$$



# Global fit results

W boson and top mass correlation – **impressive consistency of the SM**



Once  $M_H$  fixed, the SM is cornered

Effects of new physics can enter through loop corrections

→ Improve measurements of EW precision observables

# Global fit results

## Constraints on new physics

### Parametrise contributions from vacuum polarisation

- ▶ Sensitivity to new physics
- ▶ SM reference chosen to be  $M_{H,\text{ref}} = 126 \text{ GeV}$ ,  $m_{t,\text{ref}} = 173 \text{ GeV}$
- ▶  $S$ ,  $T$  depend logarithmically on  $M_H$

- ▶ Fit result:

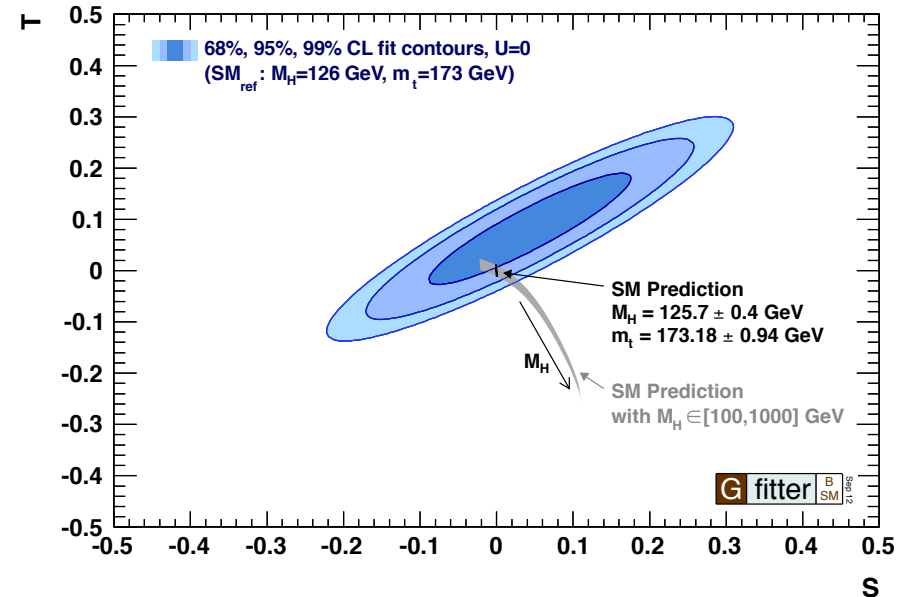
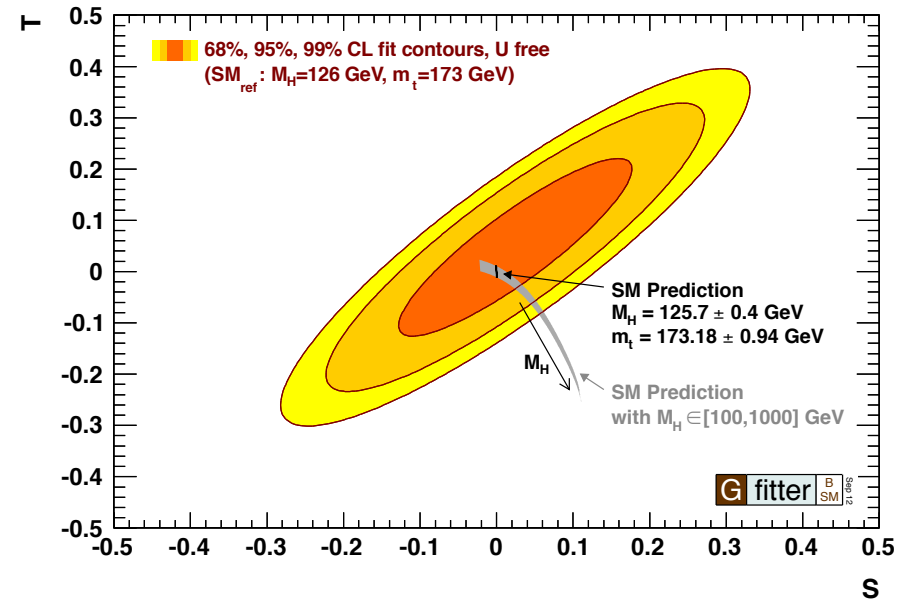
$$S = 0.03 \pm 0.10$$

$$T = 0.05 \pm 0.12$$

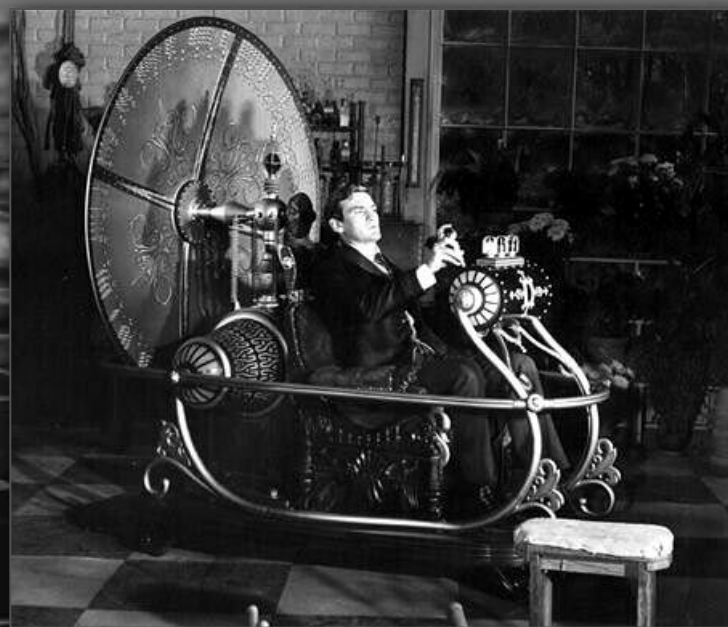
$$U = 0.03 \pm 0.10$$

with large correlation between  $S$  and  $T$

- ▶ Stronger constraints from fit with  $U = 0$
- ▶  $S$ ,  $T$ ,  $U$  fit used to constrain new physics models (Little Higgs, 2HDM, SUSY, universal extra dimensions, Technicolor, ...)



# Future



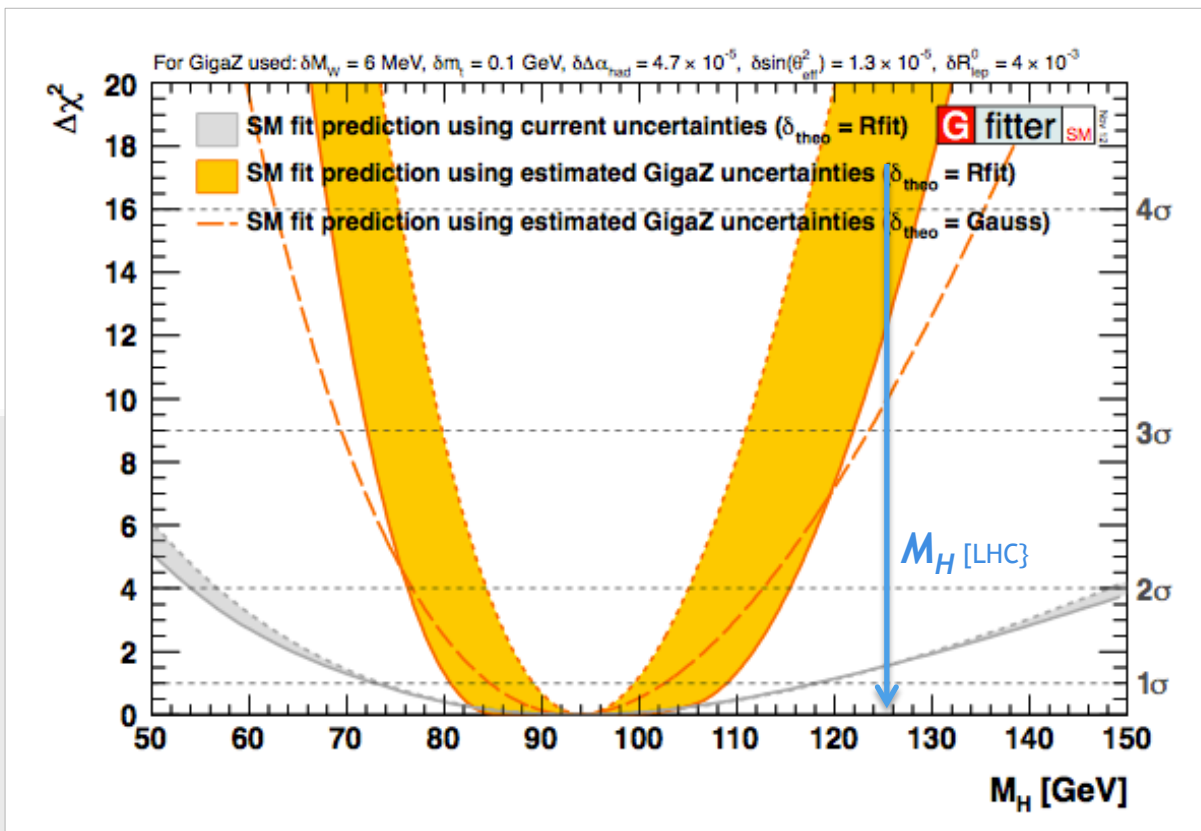
A future linear collider could tremendously improve the precision of the electroweak observables

## ILC with GigaZ

- ▶  $t\bar{t}$  threshold: obtain  $m_{\text{top}}$  from production cross section:  $\delta(m_{\text{top}}) \sim 0.1 \text{ GeV}$
- ▶ Z peak measurements
  - Polarised beams, uncertainty  $\delta A_{\text{LR}}^{0,f} : 10^{-3} \rightarrow 10^{-4}$   
translates into  $\delta \sin^2 \theta_{\text{eff}}^l : 10^{-4} \rightarrow 1.3 \times 10^{-5}$
  - High statistics:  $10^9$  Z decays:  $\delta R_l^0 : 2.5 \times 10^{-2} \rightarrow 4 \times 10^{-3}$
- ▶ WW threshold: from threshold scan:  $\delta M_W = 15 \rightarrow 6 \text{ MeV}$
- ▶ Low energy data:  $\Delta \alpha_{\text{had}}$ : more precise cross section data for low energy ( $\sqrt{s} < 1.8 \text{ GeV}$ ) and around cc resonance (BES-III), improved  $\alpha_S$ , improvements in theory:  $1.0 \times 10^{-4} \rightarrow 0.5 \times 10^{-4}$

# A future linear collider could tremendously improve the precision of the electroweak observables

Current theory uncertainties



Prospects for ILC with GigaZ

No theory uncertainty:  $M_H = 94.2^{+5.3}_{-5.0} \left[ \begin{array}{l} +23 \\ -19 \end{array} \right] \text{ GeV}$

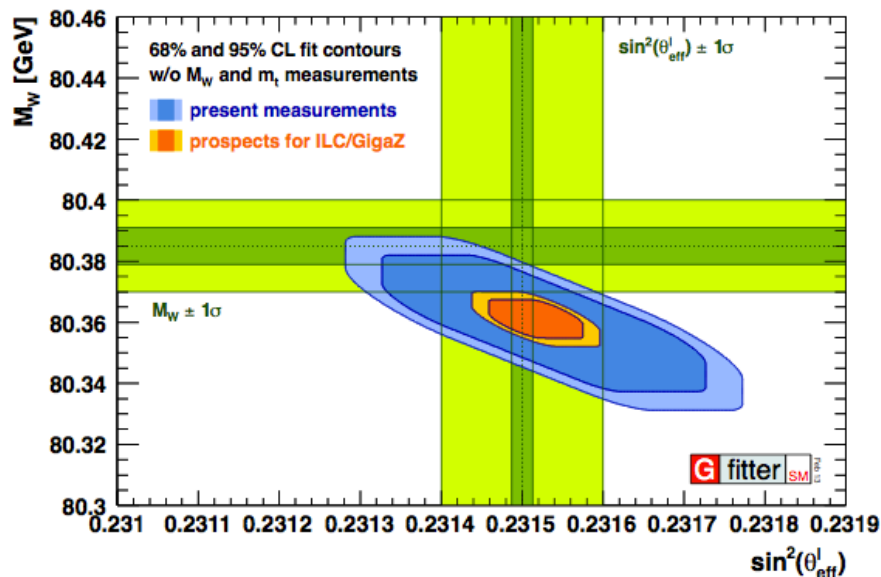
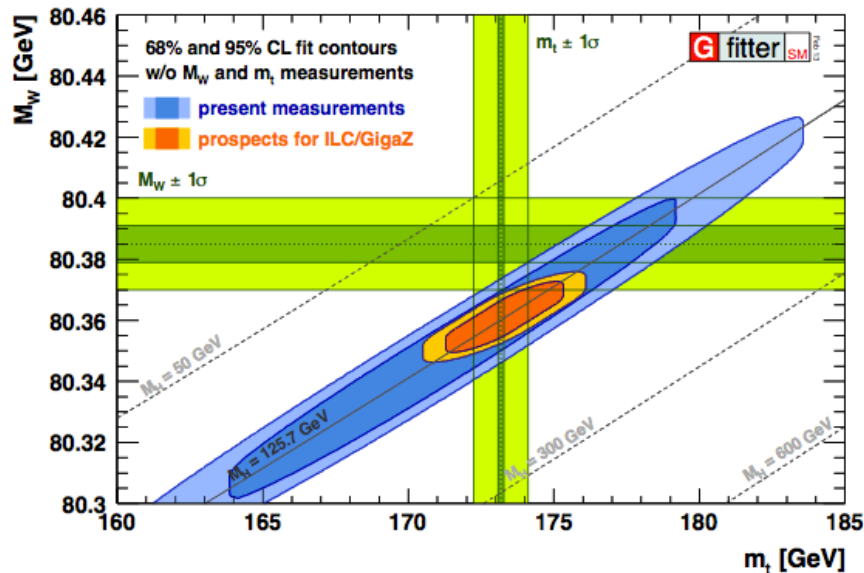
Rfit scheme:  $M_H = 92.3^{+17}_{-12} \left[ \begin{array}{l} +36 \\ -23 \end{array} \right] \text{ GeV}$

In brackets the  $4\sigma$  uncertainties

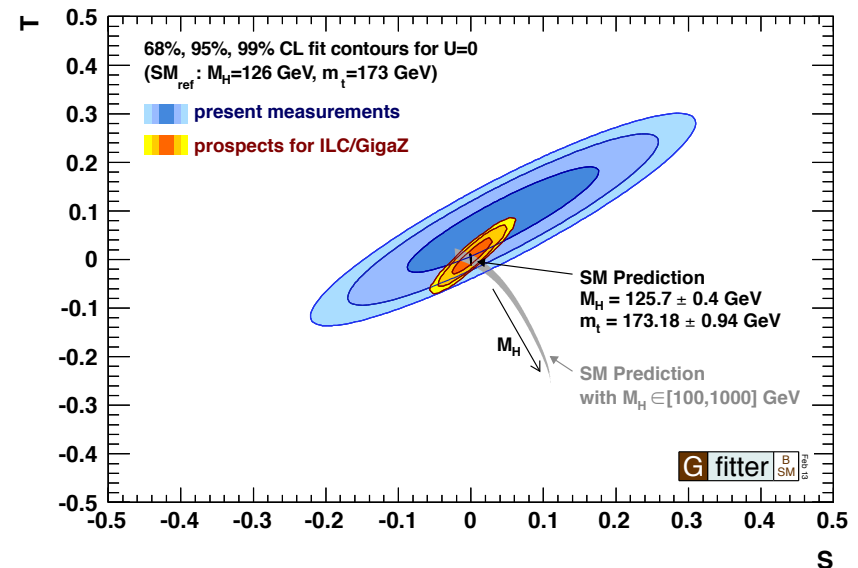


# A future linear collider could tremendously improve the precision of the electroweak observables

Prospects for ILC with GigaZ



- ▶ Assume 50% of today's theoretical uncertainty (implies three-loop EW calculations), treated *à la Rfit*
- ▶ Fit features huge uncertainty reduction for indirect determinations
- ▶ Strong constraints on  $S$ ,  $T$ ,  $U$



# Summary

Knowledge of  $M_H$  over-constrains EW fit allowing a precise prediction of observables

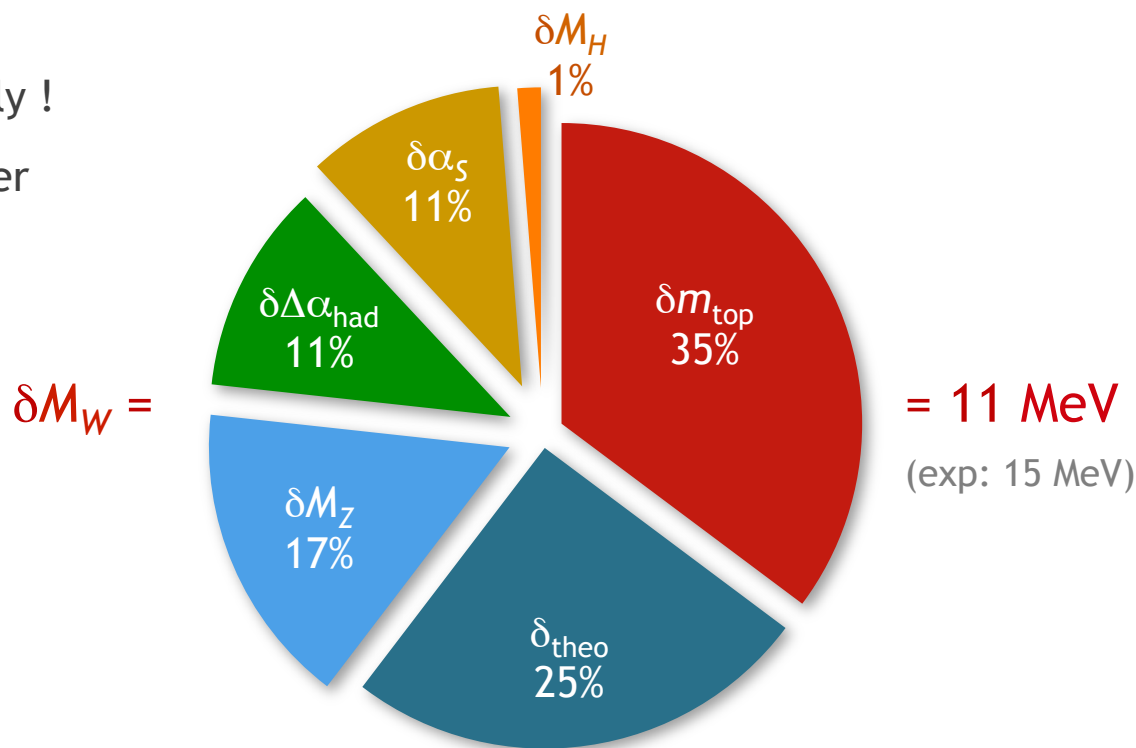
[http://gfitter.desy.de/Standard\\_Model](http://gfitter.desy.de/Standard_Model)

## SM Fit with $p$ -value of 0.07

- ▶ Incentive to revisit  $Z \rightarrow bb$  experimentally and theoretically !
- ▶ Incentive also to compute higher order contributions to other partial width observables

Significant improvement in SM prediction of key observables with  $M_H$

- ▶  $M_W$  : 28  $\rightarrow$  11 MeV
- ▶  $\sin 2\theta_{\text{eff}}^l$  : 2.3  $\rightarrow$   $1.0 \times 10^{-5}$
- ▶  $m_{\text{top}}$  : 6.2  $\rightarrow$  2.5 GeV



Improved accuracy sets benchmark for new direct measurements

Extra slides...

Parameter	Input value	Free in fit	Predicted fit result
$M_H$ [GeV]	$125.8 \pm 0.1$	yes	$125.0^{+12}_{-10}$
$M_W$ [GeV]	$80.378 \pm 0.006$	–	$80.361 \pm 0.005$
$\Gamma_W$ [GeV]	–	–	$2.0910 \pm 0.0004$
$M_Z$ [GeV]	$91.1875 \pm 0.0021$	yes	$91.1878 \pm 0.0046$
$\Gamma_Z$ [GeV]	–	–	$2.4953 \pm 0.0003$
$\sigma_{\text{had}}^0$ [nb]	–	–	$41.479 \pm 0.003$
$R_l^0$	$20.742 \pm 0.003$	–	–
$A_{\text{FB}}^{0,l}$	–	–	$0.01622 \pm 0.00002$
$A_\ell$	–	–	$0.14706 \pm 0.00010$
$\sin^2\theta_{\text{eff}}^\ell$	$0.231385 \pm 0.000013$	–	$0.23152 \pm 0.00004$
$A_c$	–	–	$0.66791 \pm 0.00005$
$A_b$	–	–	$0.93462 \pm 0.00002$
$A_{\text{FB}}^{0,c}$	–	–	$0.07367 \pm 0.00006$
$A_{\text{FB}}^{0,b}$	–	–	$0.10308 \pm 0.00007$
$R_c^0$	–	–	$0.17223 \pm 0.00001$
$R_b^0$	–	–	$0.214746 \pm 0.000004$
$\bar{m}_c$ [GeV]	$1.27^{+0.07}_{-0.11}$	yes	–
$\bar{m}_b$ [GeV]	$4.20^{+0.17}_{-0.07}$	yes	–
$m_t$ [GeV]	$173.18 \pm 0.10$	yes	$173.3 \pm 1.2$
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$ ( $\Delta$ )	$2757.0 \pm 4.7$	yes	$2757 \pm 10$
$\alpha_s(M_Z^2)$	–	yes	$0.1190 \pm 0.0005$
$\delta_{\text{th}} M_W$ [MeV]	$[-2.0, 2.0]_{\text{theo}}$	yes	–
$\delta_{\text{th}} \sin^2\theta_{\text{eff}}^\ell$ ( $\Delta$ )	$[-1.5, 1.5]_{\text{theo}}$	yes	–

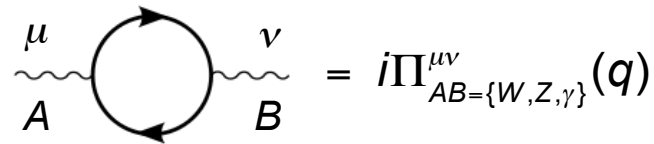
( $\Delta$ ) In units of  $10^{-5}$ . ( $\nabla$ ) Rescaled due to  $\alpha_s$  dependency.

# Oblique Corrections

Parametrising new physics contributions to electroweak precision observables

At low energies, BSM physics appears dominantly through vacuum polarisation

- ▶ Aka, oblique corrections



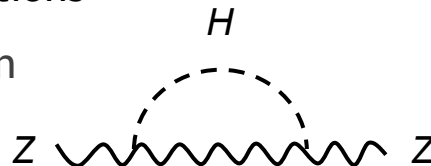
$$\text{Diagram} = i\Pi_{AB=\{W,Z,\gamma\}}^{\mu\nu}(q)$$

- ▶ Direct corrections (vertex, box, bremsstrahlung) generally suppressed by  $m_f/\Lambda$

Oblique corrections reabsorbed into electroweak parameters  $\Delta\rho$ ,  $\Delta\kappa$ ,  $\Delta r$

Electroweak fit sensitive to BSM physics through oblique corrections

- ▶ In direct competition with Higgs loop corrections



- ▶ Oblique corrections from New Physics described through “**STU parameters**”

[Peskin-Takeuchi, Phys. Rev. D46, 381 (1992)]

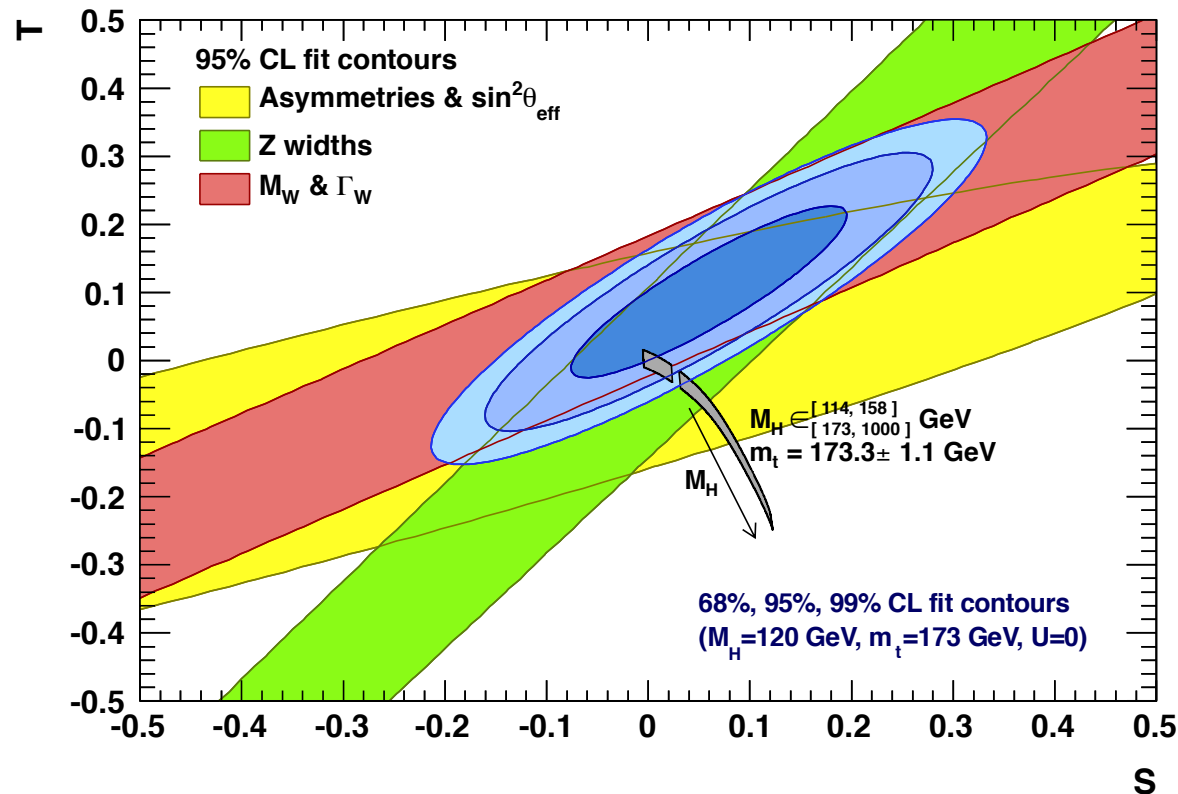
$$O_{\text{meas}} = O_{\text{SM,ref}}(M_H, m_t) + c_S S + c_T T + c_U U$$

- S** : (**S+U**) New Physics contributions to **neutral (charged) currents**
- T** : Difference between neutral and charged current processes - sensitive to **weak isospin violation**
- U** : Constrained by  $M_W$  and  $\Gamma_W$ . Usually very small in NP models (often:  $U=0$ )

# Oblique Corrections

Parametrising new physics contributions to electroweak precision observables

At low energies, BSM physics appears dominantly through vacuum polarisation



Illustrations of constraints on  $S$ ,  $T$

[ arXiv:1107.0975 ]

# Oblique Corrections

Parametrising new physics contributions to electroweak precision observables

## Definitions of $S, T, U, V, W, X$ :

[STU parameters suffice when  $(q/M)^2$  small, so that linear approximation is accurate]

$$\frac{\alpha S}{4s_W^2 c_W^2} = \left[ \frac{\delta\Pi_{ZZ}(M_Z^2) - \delta\Pi_{ZZ}(0)}{M_Z^2} \right] - \frac{(c_W^2 - s_W^2)}{s_W c_W} \delta\Pi'_{Z\gamma}(0) - \delta\Pi'_{\gamma\gamma}(0) ,$$

$$\alpha T = \frac{\delta\Pi_{WW}(0)}{M_W^2} - \frac{\delta\Pi_{ZZ}(0)}{M_Z^2} ,$$

$$\frac{\alpha U}{4s_W^2} = \left[ \frac{\delta\Pi_{WW}(M_W^2) - \delta\Pi_{WW}(0)}{M_W^2} \right] - c_W^2 \left[ \frac{\delta\Pi_{ZZ}(M_Z^2) - \delta\Pi_{ZZ}(0)}{M_Z^2} \right] - s_W^2 \delta\Pi'_{\gamma\gamma}(0) - 2s_W c_W \delta\Pi'_{Z\gamma}(0) ,$$

$$\alpha V = \delta\Pi'_{ZZ}(M_Z^2) - \left[ \frac{\delta\Pi_{ZZ}(M_Z^2) - \delta\Pi_{ZZ}(0)}{M_Z^2} \right] ,$$

$$\alpha W = \delta\Pi'_{WW}(M_W^2) - \left[ \frac{\delta\Pi_{WW}(M_W^2) - \delta\Pi_{WW}(0)}{M_W^2} \right] ,$$

$$\alpha X = -s_W c_W \left[ \frac{\delta\Pi_{Z\gamma}(M_Z^2)}{M_Z^2} - \delta\Pi'_{Z\gamma}(0) \right] .$$

[Burgess et al., PLB 326, 276 (1994), PRD 49, 6115 (1994)]

# Oblique Corrections

Parametrising new physics contributions to electroweak precision observables

Dependence of electroweak observables on  $S, T, U, V, W, X$ .

[The numerical values are based on  $\alpha^{-1}(M_Z) = 128$  and  $\sin^2\theta_W = 0.23$ ]

[Burgess et al., PLB 326, 276 (1994), PRD 49, 6115 (1994)]

$$\Gamma_Z = (\Gamma_Z)_{\text{SM}} - 0.00961S + 0.0263T + 0.0194V - 0.0207X \text{ [GeV]}$$

$$\Gamma_{bb} = (\Gamma_{bb})_{\text{SM}} - 0.00171S + 0.00416T + 0.00295V - 0.00369X \text{ [GeV]}$$

$$\Gamma_{\ell^+\ell^-} = (\Gamma_{\ell^+\ell^-})_{\text{SM}} - 0.000192S + 0.000790T + 0.000653V - 0.000416X \text{ [GeV]}$$

$$\Gamma_{\text{had}} = (\Gamma_{\text{had}})_{\text{SM}} - 0.00901S + 0.0200T + 0.0136V - 0.0195X \text{ [GeV]}$$

$$A_{\text{FB}(\mu)} = (A_{\text{FB}(\mu)})_{\text{SM}} - 0.00677S + 0.00479T - 0.0146X$$

$$A_{\text{pol}(\tau)} = (A_{\text{pol}(\tau)})_{\text{SM}} - 0.0284S + 0.0201T - 0.0613X$$

$$A_{e(P\tau)} = (A_{e(P\tau)})_{\text{SM}} - 0.0284S + 0.0201T - 0.0613X$$

$$A_{\text{FB}(b)} = (A_{\text{FB}(b)})_{\text{SM}} - 0.0188S + 0.0131T - 0.0406X$$

$$A_{\text{FB}(c)} = (A_{\text{FB}(c)})_{\text{SM}} - 0.0147S + 0.0104T - 0.03175X$$

$$A_{\text{LR}} = (A_{\text{LR}})_{\text{SM}} - 0.0284S + 0.0201T - 0.0613X$$

$$M_W^2 = (M_W^2)_{\text{SM}} (1 - 0.00723S + 0.0111T + 0.00849U)$$

$$\Gamma_W = (\Gamma_W)_{\text{SM}} (1 - 0.00723S - 0.00333T + 0.00849U + 0.00781W)$$

$$g_L^2 = (g_L^2)_{\text{SM}} - 0.00269S + 0.00663T$$

$$g_R^2 = (g_R^2)_{\text{SM}} + 0.000937S - 0.000192T$$

$$g_{V,(ve\rightarrow ve)}^e = (g_V^e)_{\text{SM}} + 0.00723S - 0.00541T$$

$$g_{A,(ve\rightarrow ve)}^e = (g_A^e)_{\text{SM}} - 0.00395T$$

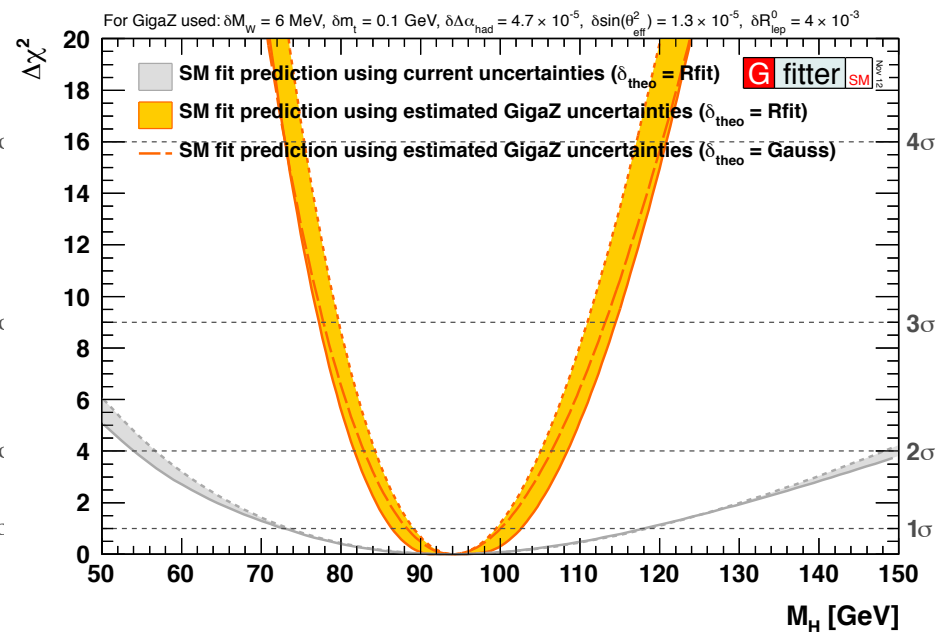
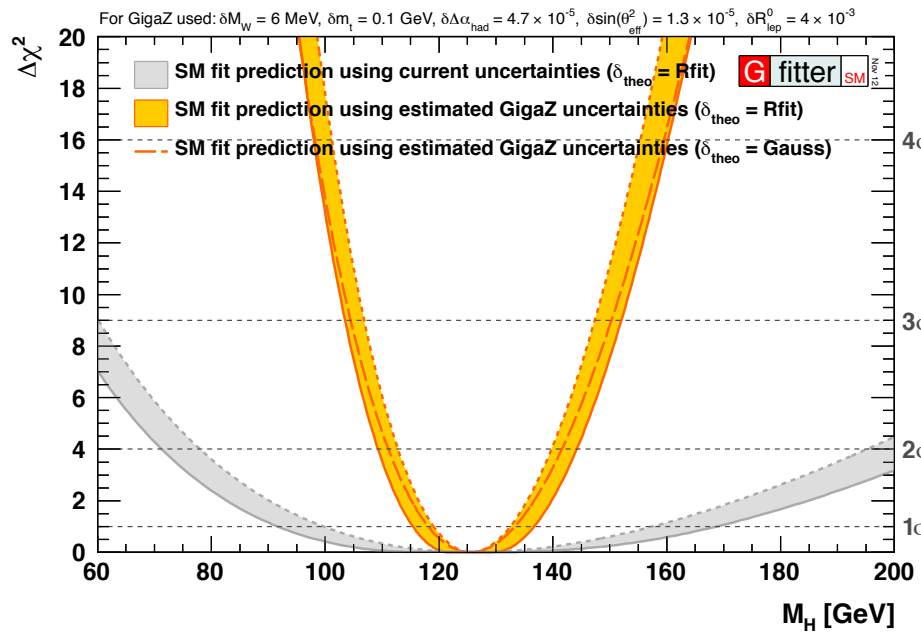
$$Q_W(^{133}\text{Cs}) = Q_W(\text{Cs})_{\text{SM}} - 0.795S - 0.0116T$$



# ILC with GigaZ

## Prospects for $M_H$ fit at ILC with GigaZ

Central values of input observables chosen to agree with their SM prediction for a Higgs mass of 126 GeV (left) and 94 GeV (right), respectively.





Status: Moriond QCD, 2013

- The branching ratio  $R_b^0$ : partial decay width of  $Z \rightarrow bb$  to  $Z \rightarrow qq$
- Freitas et al: full 2-loop calculation of  $Z \rightarrow bb$
- Contribution of same terms as in the calculation of  $\sin^2\theta_{\text{eff}}^{bb}$   
→ cross-check of two results found good agreement
- Two-loop corrections comparable to experimental uncertainty ( $6.6 \times 10^{-4}$ )

$M_H$ [GeV]	1-loop EW and QCD correction to FSR $\mathcal{O}(\alpha) + \text{FSR}_{1\text{-loop}}$ [ $10^{-3}$ ]	2-loop EW correction $\mathcal{O}(\alpha_{\text{ferm}}^2)$ [ $10^{-4}$ ]	2-loop EW and 2+3-loop QCD correction to FSR $\mathcal{O}(\alpha_{\text{ferm}}^2) + \text{FSR}_{>1\text{-loop}}$ [ $10^{-4}$ ]	1+2-loop QCD correction to gauge boson self-energies $\mathcal{O}(\alpha\alpha_s, \alpha\alpha_s^2)$ [ $10^{-4}$ ]
100	-3.632	-6.569	-9.333	-0.404
200	-3.651	-6.573	-9.332	-0.404
400	-3.675	-6.581	-9.331	-0.404

# Higgs couplings in the EW fit



Status: Moriond QCD, 2013

- In latest ATLAS  $H \rightarrow \gamma\gamma$ ,  $2.3\sigma$  deviation seen from SM  $\mu (\equiv 1.0)$
- Interpret.:  $H \rightarrow VV$  couplings scaled with  $c_V$

From: Falkowski et al, arXiv:1303.1812

- Modified Higgs couplings can be constrained by EW fit through extended STU formalism.
- Result of  $c_V$  driven by limit on  $T$  parameter.
  - Tree-level relation:  $\rho_0 = \frac{M_{W_0}^2}{M_{Z_0}^2 c_W^2} = 1 + \alpha T$
  - $\alpha T \approx \frac{3g_Y^2}{32\pi^2} (c_V^2 - 1) \log(\Lambda/m_Z)$
  - Reminder:  $T = 0.05 \pm 0.12$  (Gfitter)
- EW-fit Falkowski et al:  $c_V \approx 1.08 \pm 0.07$ 
  - Blue dashed:  $c_V$  from  $\mu$ 's, black: comb. w/ EW

